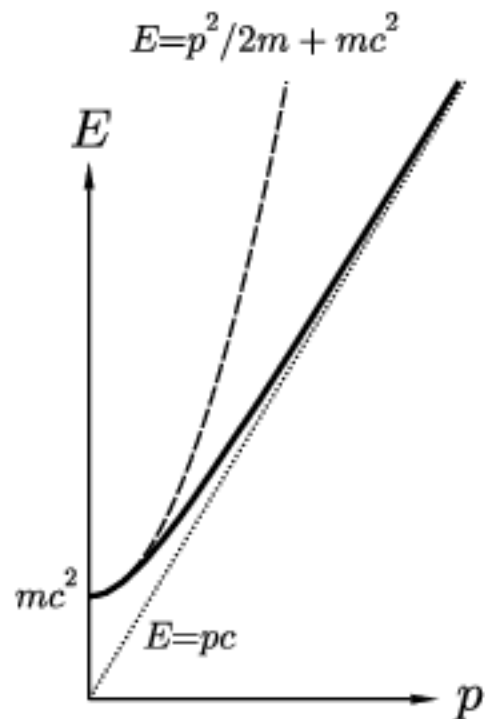

Lecture 25. *Relativistic Gas*

- In this lecture, we will focus on cases where classical mechanics description of kinetic energy fails at high velocity

Dispersion relation



Kinetic Energy, relativistic formula

$$E^2 = p^2 c^2 + m^2 c^4$$

m : rest mass

Non-relativistic limit

$$p/m \ll c \longrightarrow E = \frac{p^2}{2m} + mc^2$$

mc^2 : rest mass energy

$$E(k) = \hbar^2 k^2 / 2m$$

Ultra-relativistic limit

$$p/m \gg c \longrightarrow E = pc$$

e.g., photons

$$\omega = ck$$

Ultrarelativistic gas $E = pc$

We will consider an UR gas of finite mass particles

Partition function of a single-particle: $Z_1 = \int_0^\infty e^{-\beta\hbar kc} g(k) dk$

Density of states: $g(k) dk = \frac{V k^2 dk}{2\pi^2}$

$$\longrightarrow Z_1 = \frac{V}{2\pi^2} \left(\frac{1}{\beta\hbar c} \right)^3 \int_0^\infty e^{-x} x^2 dx \quad x = \beta\hbar kc.$$

$$\longrightarrow Z_1 = \frac{V}{\pi^2} \left(\frac{k_B T}{\hbar c} \right)^3$$

$$Z_1 = \frac{V}{\Lambda^3} \quad \Lambda = \frac{c\pi^{2/3}}{k_B T}$$

What changed? T^{-3} instead of $T^{-3/2}$

Thermodynamic properties!

- Partition function of the ultra relativistic gas of *indistinguishable* particles

$$Z_N = \frac{Z_1^N}{N!} \quad \text{*note: when is this valid?}$$

$$\ln Z_N = N \ln V + 3N \ln T + \text{constants}$$

- Internal energy $U = -\frac{d \ln Z_N}{d\beta} = 3Nk_B T$

- Heat capacity $C_V = \left(\frac{\partial U}{\partial T}\right)_V \longrightarrow C_V = 3Nk_B$

*note: why is this different from the equipartition theorem?

$$\ln Z_N = N \ln V + 3N \ln T + \text{constants}$$

- **Internal energy** $U = -\frac{d \ln Z_N}{d\beta} = 3Nk_B T$

- **Heat capacity** $C_V = \left(\frac{\partial U}{\partial T}\right)_V \longrightarrow C_V = 3Nk_B$

- **Helmholtz function** $F = -k_B T \ln Z_N = -k_B T N \ln V - 3Nk_B T \ln T - k_B T \times \text{constants}$

- **Pressure** $p = -\left(\frac{\partial F}{\partial V}\right)_T = \frac{Nk_B T}{V} = nk_B T$ *ideal gas law!

- **Enthalpy** $H = U + pV = 4Nk_B T$

Entropy

$$\Lambda = \frac{c\pi^{2/3}}{k_B T}$$

$$\begin{aligned}\ln Z_N &= N \ln V - 3N \ln \Lambda - N \ln N + N \\ &= N \ln \left(\frac{1}{n\Lambda^3} \right) + N\end{aligned}$$

$$\begin{aligned}F &= -k_B T \ln Z_N \\ &= Nk_B T [\ln(n\Lambda^3) - 1], \\ S &= \frac{U - F}{T} \\ &= Nk_B [4 - \ln(n\Lambda^3)], \\ G &= H - TS = 4Nk_B T - Nk_B T [4 - \ln(n\Lambda^3)] \\ &= Nk_B T \ln(n\Lambda^3).\end{aligned}$$

Pressure

$$p = \frac{u}{3}$$

$$p = \frac{2u}{3}$$

Adiabatic Expansion of an ultra relativistic gas

- No heat in or out!

- No entropy change $S = Nk_B[4 - \ln(n\Lambda^3)]$

- Thus this is also a constant $n\Lambda^3 = n\left(\frac{\hbar c\pi^{2/3}}{k_B T}\right)^3 = \frac{N}{V}\left(\frac{\hbar c\pi^{2/3}}{k_B T}\right)^3$

Expansion of the universe

$$VT^3 = \text{constant}$$

$$T \propto V^{-1/3} \propto a^{-1} \quad \bullet \quad \alpha \text{ is a scale factor}$$

- The temperature of the CMB is inversely proportional to the scale factor of the universe.

- Non-relativistic gas: $VT^{3/2} = \text{constant}$ $T \propto V^{-2/3} \propto a^{-2}$,

- **Faster cooling than UR**

Density of the expanding universe

- Non-relativistic gas: $\rho \propto V^{-1} \longrightarrow \rho \propto a^{-3}$
- Relativistic gas: $\rho = \frac{u}{c^2} \xrightarrow{u = 3p} \rho \propto a^{-4}$
 $p \propto V^{-4/3}$

The density of relativistic particles decreases faster than that of non-relativistic particles, as the universe expands and temperature goes down

Early universe was radiation dominated and the Universe has become matter dominated

Summary

| Property | Non-relativistic | Ultrarelativistic |
|---------------------|---|---|
| Z_1 | $\frac{V}{\lambda_{\text{th}}^3}$ $\lambda_{\text{th}} = \frac{h}{\sqrt{2\pi mk_{\text{B}}T}}$ | $\frac{V}{\Lambda^3}$ $\Lambda = \frac{c\pi^{2/3}}{k_{\text{B}}T}$ |
| U | $\frac{3}{2}Nk_{\text{B}}T$ | $3Nk_{\text{B}}T$ |
| H | $\frac{5}{2}Nk_{\text{B}}T$ | $4Nk_{\text{B}}T$ |
| p | $\frac{Nk_{\text{B}}T}{V}$ $= \frac{2n}{3}$ | $\frac{Nk_{\text{B}}T}{V}$ $= \frac{n}{3}$ |
| F | $Nk_{\text{B}}T[\ln(n\lambda_{\text{th}}^3) - 1]$ | $Nk_{\text{B}}T[\ln(n\Lambda^3) - 1]$ |
| S | $Nk_{\text{B}}[\frac{5}{2} - \ln(n\lambda_{\text{th}}^3)]$ | $Nk_{\text{B}}[4 - \ln(n\Lambda^3)]$ |
| G | $Nk_{\text{B}}T \ln(n\lambda_{\text{th}}^3)$ | $Nk_{\text{B}}T \ln(n\Lambda^3)$ |
| Adiabatic expansion | $VT^{3/2} = \text{constant}$ $pV^{5/3} = \text{constant}$ | $VT^3 = \text{constant}$ $pV^{4/3} = \text{constant}$ |