
Lecture 27. *Cooling Real Gas*

- In this lecture, we will move beyond the ideal gas description and revisit fundamental aspects like Joule expansion
- We will also discuss the *Joule-Kelvin* throttling process and how real gas can be *liquefied*

The Joule expansion (again)

Irreversible process

Reminder:

- (1) We have an isolated system (no heat in or out)
- (2) No work done either

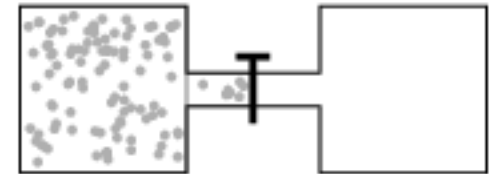
Effect of intermolecular interactions ?

Does the gas warm, cool, or remain at constant T?

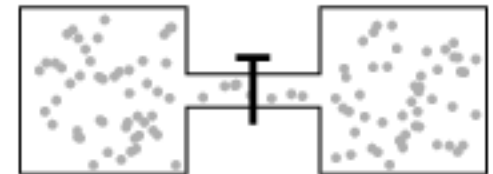
$$\mu_J = \left(\frac{\partial T}{\partial V} \right)_U$$

Joule Coefficient

Initial



Final



$$\mu_J = - \left(\frac{\partial T}{\partial U} \right)_V \left(\frac{\partial U}{\partial V} \right)_T$$

$$= - \frac{1}{C_V} \left(\frac{\partial U}{\partial V} \right)_T$$

$$\left(\frac{\partial x}{\partial y} \right)_z = - \left(\frac{\partial x}{\partial z} \right)_y \left(\frac{\partial z}{\partial y} \right)_x$$

$$dU = TdS - pdV \rightarrow \left(\frac{\partial U}{\partial V} \right)_T = T \left(\frac{\partial S}{\partial V} \right)_T - p$$

Maxwell relation (*)

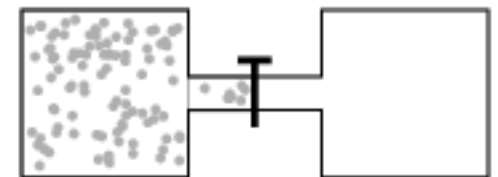
$$\left(\frac{\partial U}{\partial V} \right)_T = T \left(\frac{\partial p}{\partial T} \right)_V - p$$

$$\mu_J = - \frac{1}{C_V} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right]$$

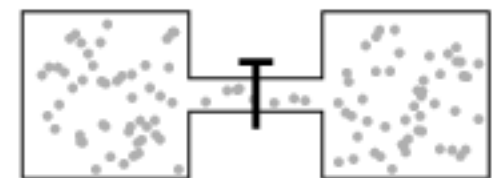
$$\mu_J = \left(\frac{\partial T}{\partial V} \right)_U$$

Joule Coefficient

Initial



Final



$$\mu_J = -\frac{1}{C_V} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right]$$

Ideal gas: $p = RT/V$
 $(\partial p / \partial T)_V = R/V$

$$\mu_J = 0$$

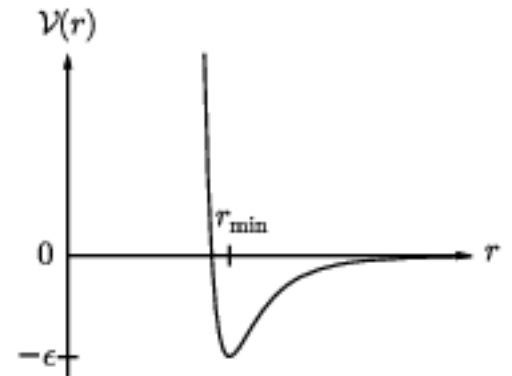
Physical Interpretation of cooling

When a gas freely expands into a vacuum, the time-averaged distance between neighboring molecules increases

Thus, the magnitude of the potential energy resulting from the attractive intermolecular interactions is reduced.

This leads to an increase in total energy

Since U is constant in a Joule expansion, the kinetic energy must be reduced and hence the temperature falls.



van der Waals gas

$$\mu_J = -\frac{1}{C_V} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right]$$

$$p = RT/(V - b) - a/V^2$$

$$\left(\frac{\partial p}{\partial T} \right)_V = \frac{R}{V - b}$$

$$\mu_J = -\frac{1}{C_V} \left[\frac{RT}{V - b} - \frac{RT}{V - b} + \frac{a}{V^2} \right] = -\frac{a}{C_V V^2}$$

Temperature change:

$$\Delta T = -\frac{a}{C_V} \int_{V_1}^{V_2} \frac{dV}{V^2} = -\frac{a}{C_V} \left(\frac{1}{V_1} - \frac{1}{V_2} \right) < 0$$

Isothermal expansion

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial p}{\partial T}\right)_V - p \longrightarrow \Delta U = \int_{V_1}^{V_2} \left[T \left(\frac{\partial p}{\partial T}\right)_V - p \right] dV.$$

Ideal gas: no change!

vdW gas: $\Delta U = \int_{V_1}^{V_2} \frac{a}{V^2} dV = a(1/V_1 - 1/V_2)$ (no b dependence!)

Entropy of vdW gas

$$dS = \left(\frac{\partial S}{\partial T}\right)_V dT + \left(\frac{\partial S}{\partial V}\right)_T dV$$

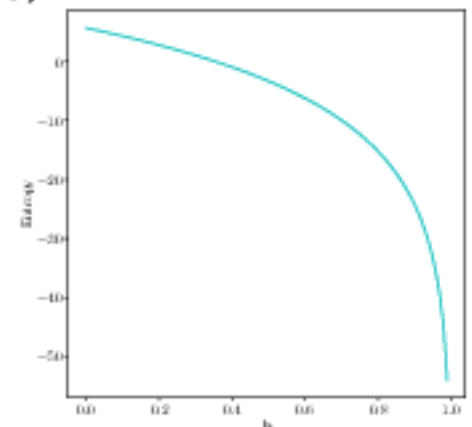
$$= \frac{C_V}{T} dT + \left(\frac{\partial p}{\partial T}\right)_V dV \quad (\text{See chapter 16})$$

$$\left(\frac{\partial p}{\partial T}\right)_V = R/(V - b)$$



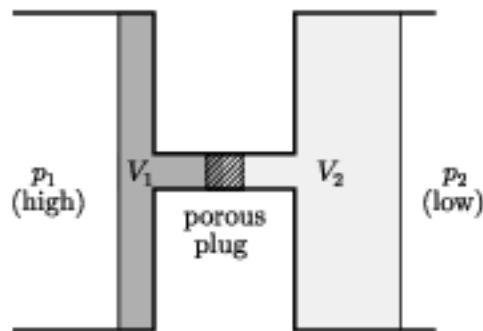
$$S = C_V \ln T + R \ln(V - b) + \text{constant}$$

$V = 1, T = 100, C_V = 1$



Only depends on b , not a , as b governs the number of microstates available.

Joule-Kelvin Expansion



Change in internal energy of gas moving from left to right?

- (1) work done to move V_1 at p_1 : $p_1 V_1$
- (2) Work to move V_2 at p_2 : $p_2 V_2$
- (3) Change in internal energy is the difference between the work done on the left and on the right

$$U_1 + p_1 V_1 = U_2 + p_2 V_2$$

$$\longrightarrow H_1 = H_2 \quad (\text{enthalpy is conserved})$$

Temperature change? $\mu_{JK} = \left(\frac{\partial T}{\partial p} \right)_H$

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$$\left(\frac{\partial x}{\partial y} \right)_z = - \left(\frac{\partial x}{\partial z} \right)_y \left(\frac{\partial z}{\partial y} \right)_x$$



$$\mu_{JK} = - \left(\frac{\partial T}{\partial H} \right)_p \left(\frac{\partial H}{\partial p} \right)_T = - \frac{1}{C_p} \left(\frac{\partial H}{\partial p} \right)_T$$

$$dH = TdS + Vdp \longrightarrow \left(\frac{\partial H}{\partial p} \right)_T = T \left(\frac{\partial S}{\partial p} \right)_T + V$$

$$\longrightarrow \left(\frac{\partial H}{\partial p} \right)_T = -T \left(\frac{\partial V}{\partial T} \right)_p + V$$

Maxwell relation (*)



$$\mu_{JK} = \frac{1}{C_p} \left[T \left(\frac{\partial V}{\partial T} \right)_p - V \right]$$

Temperature change: $\Delta T = \int_{p_1}^{p_2} \frac{1}{C_p} \left[T \left(\frac{\partial V}{\partial T} \right)_p - V \right] dp$

$$\mu_J = - \frac{1}{C_V} \left[T \left(\frac{\partial p}{\partial T} \right)_V - p \right]$$

Entropy Change

$$dH = TdS + Vdp = 0 \quad \longrightarrow \quad \Delta S = - \int_{p_1}^{p_2} \frac{V}{T} dp$$

Ideal gas: $R \ln(p_1/p_2) > 0$ (irreversible process)

Heating or cooling?

- Changes in sign when $\left(\frac{\partial V}{\partial T}\right)_p = \frac{V}{T}$

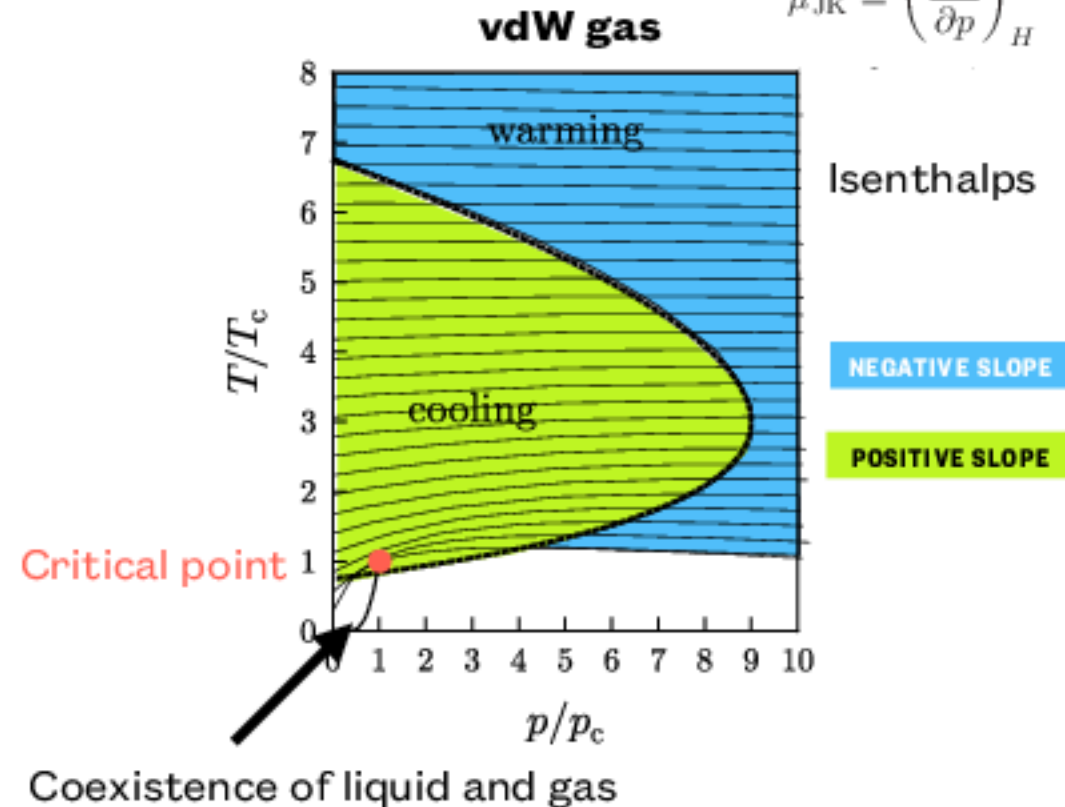
“Inversion curve” in the T-p plane

Maximum inversion temperature: under that T, JK can result in cooling

⁴ He	H ₂	N ₂	Ar	CO ₂
43	204	607	794	1275

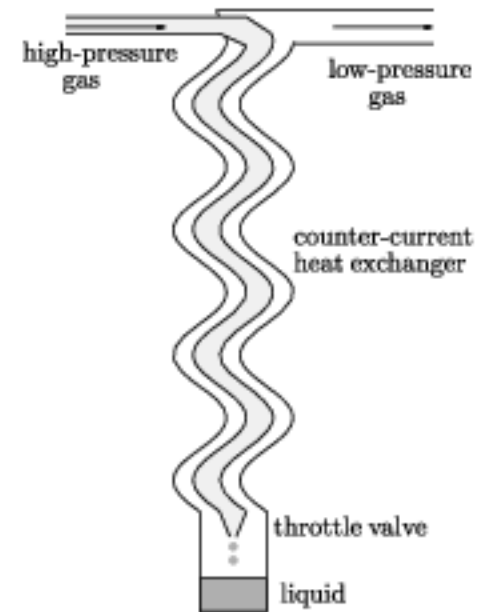
$$\mu_{JK} = \frac{1}{C_p} \left[T \left(\frac{\partial V}{\partial T} \right)_p - V \right]$$

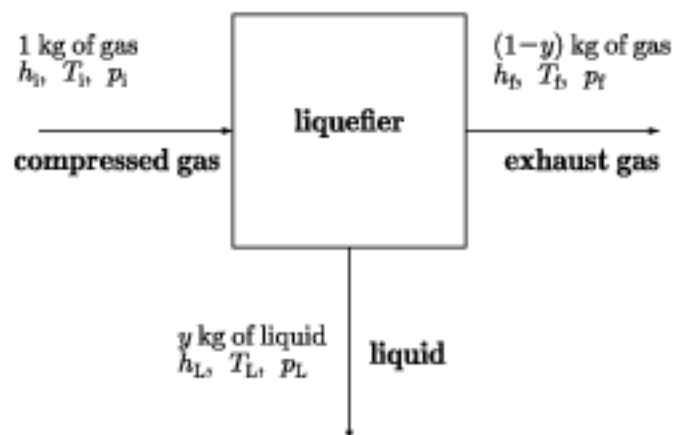
$$\mu_{JK} = \left(\frac{\partial T}{\partial p} \right)_H$$



Liquefaction of gases

- Can use the JK process, but must start below the maximum inversion temperature of that gas
- High-p gas is forced through a throttle valve, and results in cooling by the JK process
- This results in low-p gas plus liquid
- Counter current heat exchanger where the outgoing cold low-pressure gas is used to precool the incoming warm high-pressure gas (and stay under the maximum inversion temperature)





Enthalpy is conserved, thus:

$$h_i = yh_L + (1 - y)h_f.$$

Efficiency:

$$y = \frac{h_f - h_i}{h_f - h_L}$$

For an efficient heat exchanger, T_i (compressed gas) and T_f (exhaust gas) will be the same
 $p_f=1$ atm, T_L is fixed (because it is in equilibrium with its vapor)

Thus h_f and h_L are fixed, we can only vary h_i and maximum y is found for:

$$\left(\frac{\partial h_i}{\partial p_i}\right)_{T_i} = 0$$

$$\left(\frac{\partial h}{\partial p}\right)_T = -C_p \mu_{JK} \longrightarrow \mu_{JK} = 0.$$

We get highest efficiency on the inversion curve!

Summary

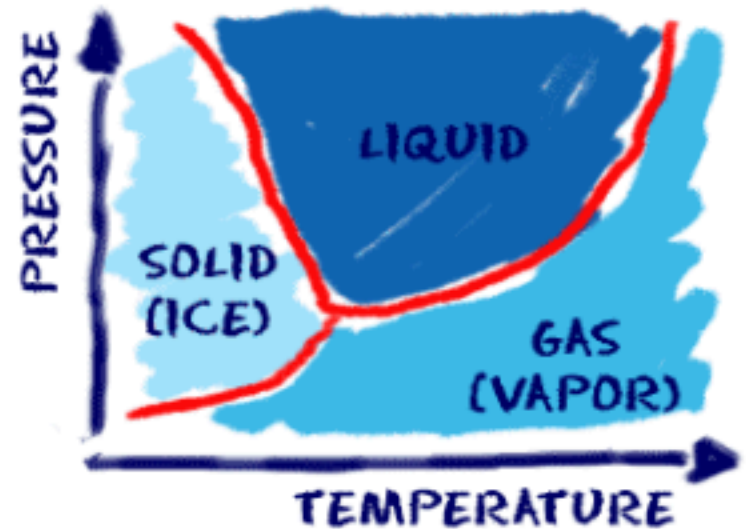
Joule expansion results in cooling for non-ideal gases because of the attractive interactions between molecules.

The entropy of a gas depends on the non-zero size of molecules.

The Joule-Kelvin expansion is a steady flow process in which enthalpy is conserved. It can result in either warming or cooling of a gas. It forms the basis of many gas liquefaction techniques.

Lecture 28. *Phase transitions*

- In this lecture, we will dig deeper into phase transitions and will study how the thermodynamics of a phase transition can be characterized using various criteria
- Exceptionally, this lecture will be provided in two parts with the second part of the screencast devoted to the so-called *Ising model*, an example of a simple system to study phase transition.



http://www.chem4kids.com/files/matter_changes.html